

# Modelling CO<sub>2</sub> - water mixture thermodynamics using various equations of state (EoSs) with emphasis on the potential of the SPUNG EoS

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## Abstract

CO<sub>2</sub> -water is a very important mixture in the Carbon Capture and Storage (CCS) industry. The mixture can have a broad range of concentrations, from water as an impurity in CO<sub>2</sub> transport to high water concentrations in sequestration processes. CO<sub>2</sub> -water mixture is challenging due to the polar nature that induces difficulties describing the interaction between CO<sub>2</sub> and water when modelling the behavior. The work focus on the evaluation of the predictability of the extended corresponding state equation SPUNG in dealing with CO<sub>2</sub> -water thermodynamics. The evaluation is done by comparing the behavior of SPUNG equation of state (EoS) to experimental data, and two other EoSs of a different class. The two other EoSs are the cubic equation Soave-Redlich-Kwong (SRK) with van der Waals mixing rules, and SRK with Huron-Vidal mixing rules.

The predictability of the single and liquid rich phases densities, two-phase solubilities and dew line are investigated over a wide range of pressures, temperatures and mixture compositions. The results show better density prediction using SPUNG EoS over all the evaluated conditions compared to SRKs with a potential of improvements by changing the reference fluid. However, the CO<sub>2</sub> solubility prediction using SPUNG requires the use of other mixing rules that can account for the polar nature of the system.

*Keywords:* Cubic EoS, Extended Corresponding States, MBWR EoS, CCS, Reference fluid, VLE.

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## 1. Introduction

Through the various CCS processes, CO<sub>2</sub> exists in mixtures with various impurities like CH<sub>4</sub>, CO, H<sub>2</sub>O, H<sub>2</sub>S, N<sub>2</sub>, NO<sub>2</sub> and O<sub>2</sub>. Therefore, the knowledge of the thermophysical properties of those mixtures is a key challenge for accurate design of efficient and secure processes. Hendriks et al. (2010) pointed out the need for accurate thermophysical properties.

Even in the cases where experimental data exist for a mixture, they are discrete and local in nature, and more continuous and generic solutions are rather practical. Consequently, the modelling of the thermodynamic properties for pure CO<sub>2</sub> and CO<sub>2</sub> mixtures is a very important aspect for the analysis and detailed simulation of CCS processes. Indeed, the choice of models may have a great impact on the decisions about process design, energy efficiency, economy and safety.

A computationally cheap modelling strategy is to empirically fit experimental data. This solution has a poor generality to different mixtures and for different phases and intervals outside the fitted range. A more physically profound and rather general and continuous approach is the use of Equations of State (EoSs), which will be discussed more in detail in the following section. There is a large variety of EoSs at various levels of sophistication. Cubic EoSs like Soave-Redlich-Kwong (SRK) (Soave, 1972), SRK with Huron Vidal mixing rules (SRK-HV) (Huron and Vidal, 1979) and Peng–Robinson (PR) (Peng and Robinson, 1976) are among the simplest. Multi-parameter approaches like Span–Wagner (Span and Wagner, 1996) for pure CO<sub>2</sub> and GERG (Groupe Européen de Recherches Gazières) (Kunz et al., 2007) for mixtures are at least one order of magnitude higher in computational time. The full methods of extended corresponding states like those implemented in the REFPROP library of the National Institute of Standards and Technology (NIST) (Lemmon et al., 2010) are even more expensive than the multi-parameter approaches. Among the state-of-the-art approaches are the Cubic-Plus-Association (CPA) (Kontogeorgis et al., 1996) and the Statistical Associating Fluid Theory (SAFT) (Chapman et al., 1990) EoSs. Tsvintzelis et al. (2011) and Diamantonis and Economou (2012) demonstrated the success of the CPA and Perturbed Chain SAFT (PC-SAFT), respectively, in modelling the polar mixtures including CO<sub>2</sub>-water. The two mentioned articles included reviews of the development of CPA and SAFT, and their different modifications and combinations.

The level of sophistication and generality usually has a direct relation to accuracy and computational complexity and, consequently, a trade-off arises. While the accuracy of a model is of higher importance than the computational efficiency for the process analysis, the computational complexity has a significant effect on the cost and feasibility of a CFD simulation. Three other dimensions of the challenge of developing or selecting a model are the generality with respect to different fluids and mixtures, consistency, and numerical stability when using it in conjunction with CFD simulations.

A consistent approach that is not well known but has shown a very good compromise in accuracy and computation time for hydrocarbons is the SPUNG

EoS (Jørstad, 1993). The SPUNG EoS was not published outside Jørstad thesis before the work of (Wilhelmsen et al., 2012). The latter demonstrated that SPUNG is a very good compromise for CO<sub>2</sub> with some non-polar binary and ternary impurities. They showed that for the calculations of density, enthalpy and entropy over a 10 000 random conditions in different phase regions, and for three component CO<sub>2</sub> mixtures, SPUNG run time was 4 times and GERG was 40 times of that of SRK. The work also showed that the SPUNG EoS accuracy was generally high and close to GERG and Span-Wagner for pure species as references. These results of computational time were only for single phase including critical and near critical conditions, since flash results would strongly depend on algorithmic and implementation.

Since generality is a critical aspect when selecting a model to be used for CFD, this work aims to study the behavior of the SPUNG EoS for the polar mixture of CO<sub>2</sub> and water. The study of the SPUNG EoS generality with respect to CO<sub>2</sub>-water mixtures is of particular importance because these mixtures exists commonly in the range of processes in CCS industry. Moreover, they are very challenging mixtures due to the polar nature. A preliminary study was conducted by Ibrahim et al. (2012) that covered a few conditions that exist in CCS. The study presented here is to extend the evaluation over a wide range of conditions, compositions, temperatures, pressures that might occur in various CCS processes. Consequently, this study can be used as a comprehensive visualized analysis of the behavior and the errors of each EoS at this wide range of conditions.

The study aims mainly to investigate the shortcoming and capabilities of SPUNG EoS when dealing with CO<sub>2</sub>-water. Furthermore, to investigate the impact of the EoS used for the shape factors scaling, and the choice of the reference fluid on the predictability and limitations of the method. In addition, the study aims to investigate the potential and possibilities of improvement for SPUNG EoSs.

Here, an evaluation is done by comparing the behavior of SPUNG to two other EoSs and with experimental data. The two EoSs are the cubic equations SRK with the van der Waals mixing rule (Soave, 1972) and SRK-HV (Huron and Vidal, 1979).

The SRK-HV was used because it showed very good results for the solubility predictions for the investigated mixtures as reported by Austegard et al. (2006). Furthermore, The classical SRK was chosen because SPUNG use it for computing the shape factors and because it is a simple model and is commonly used in industry.

In this work, the predictability of single phase densities, dew lines, mixture solubilities in two-phase, and rich densities will be evaluated.

## 2. Theory

### 2.1. Equations of state

An EoS is a model that calculates for both the liquid and gas phase using the same expression. This enhances the continuity near the critical point. An

EoS for an  $N_c$  component mixture can be regarded as an expression for pressure  $P$  as a function of the mole fractions  $x_i$ , the temperature  $T$  and the volume  $V$ . Given this expression, it can be manipulated to calculate the fugacity of each component. In the following subsections, a brief description of the equations of state used in the work will be given together with references for further discussions.

## 2.2. The SRKs EoSs

The classical SRK model used here (Soave, 1972) is a cubic EoS that uses van der Waals mixing rules.

The SRK-HV model, proposed by Huron and Vidal (1979), is an improvement from the classical SRK, as it derives a definition for the mixing rules from the excess Gibbs energy at infinite pressure. A detailed description of that model has been given by Solbraa (2002). The SRK-HV implementation used in this work has parameters regressed over a wide range of CO<sub>2</sub>-water data, and the regression work is described in detail by Austegard et al. (2006). The SRK-HV evaluated in this work uses the Twu-Bluck-Cunningham (TBC) (Twu et al., 1991) formulation for computing the alpha parameter.

## 2.3. The corresponding states principle

The principle of corresponding states assumes that all substances exhibit the same behavior at a reduced state. A corresponding state EoS typically has one or more reference components described very accurately by a reference EoS. Therefore, the compressibility of the investigated fluids or mixtures can be evaluated as  $Z = Z(V_{\text{Ref}}, T_{\text{Ref}}, \omega, \dots)$ . In the corresponding states approach, the reference fluid volume  $V_{\text{Ref}}$  and temperature  $T_{\text{Ref}}$  are the reduced volume and temperature,  $V_{\text{R}}$  and  $T_{\text{R}}$ , of the fluid or the mixture investigated.

## 2.4. The extended corresponding states principle

### 2.4.1. Basic concept

In the extended corresponding states concept, the mapping between the investigated fluid or mixture  $T$  and  $V$  and the reference fluid  $V_{\text{Ref}}$  and  $T_{\text{Ref}}$  is done via the shape factors  $\theta$  and  $\phi$ . These shape factors take into account how the fluids or the mixture in consideration differ from the reference fluid, where  $T_{\text{Ref}} = T/\theta$  and  $V_{\text{Ref}} = V/\phi$ . The shape factors  $\theta$  and  $\phi$  can be computed via shape factor functions, using semi-empirical functions, an accurate reference equation for each component, or using a simpler EoS. The work on shape factor functions started as early as 1968 by Leach et al. (1968). Subsequently, many contributions were made, examples are the work by Fisher and Leland (1970) and of Ely (1990), who has introduced the first exact shape factor concept. One of the most recent work on shape factor functions was conducted by Estela-Urbe and Trusler (1998). The computation of exact shape functions is computationally very expensive, which is why the concept was left behind and thought to be impractical for use with numerical simulations. However, several implementations of the concept of extended corresponding states use simpler

equations of state to compute shape factors, which show a good compromise between accuracy and computation time.

#### *2.4.2. The SPUNG EoS*

The SPUNG EoS investigated here is an instance of the extended corresponding states approach, which was enlightened by the work of Møllerup (1980) and developed first for hydrocarbons. The SPUNG EoS uses the cubic SRK EoS to calculate the shape factors and propane as a reference fluid. Furthermore, it uses the accurate modified Benedict-Webb-Rubin (MBWR) (Younglove and Ely, 1987) EoS for the reference fluid. The SPUNG EoS is described more in detail in the doctoral thesis by Jørstad (1993). It was developed for low temperature hydrocarbon mixtures, and it has improved density and enthalpy prediction while maintaining a good compromise in computational expenses. Propane was chosen as the reference fluid to ensure that the reduced temperature of the considered mixtures would always be above the reduced triple point of the reference fluid in order to avoid extrapolation of the reference equation. For CO<sub>2</sub> mixtures the choice of a different reference fluid and equation should be considered, but in this work only the original SPUNG EoS formulation was used.

### **3. Methodology**

#### *3.1. Numerical Tools*

The NTNU-SINTEF in-house thermodynamic library was mainly used for the study presented. The library is a tool for predicting the thermodynamic properties using various approaches that ranges in level of sophistication and underlying theory. The SRK-HV model used here was the one described by Austegard et al. (2006) and with the regressed coefficients listed there. The library uses a tolerance of  $10^{-4}$  for both the multi-phase flash algorithm and the compressibility factor calculations.

#### *3.2. Setup of the investigation*

##### *3.2.1. Single phase density predictions, low to moderate pressures*

The three EoSs were evaluated at a set of low to moderate pressures (up to 100 bars) that ensured a single phase at given temperatures and water concentrations in the mixture. The pressures, levels of temperature and water concentrations were chosen to enable comparison with the experimental work of Patel and Eubank (1988). Four concentrations of 98, 90, 75 and 50% H<sub>2</sub>O were evaluated. The concentration of 98% was evaluated at the temperatures of 225, 200, 100 and 50 °C. The 90% was evaluated at 200, 100 and 71 °C, the 75% was evaluated at 225, 200, and 100 °C, and finally the 50% was evaluated at 225, 200 and 125 °C. It was clear that, as the H<sub>2</sub>O concentration increased in the mixture, it was not possible to go to some low temperatures while maintaining the mixture in gaseous phase conditions. This explains the differences in the lower bound of the evaluated temperatures at the four studied concentrations.

### 3.2.2. *Single phase density predictions, high pressures*

A more challenging set of conditions at elevated pressures over a wider range of concentrations was evaluated. The set of concentrations ranged from CO<sub>2</sub> dominant (90%) to H<sub>2</sub>O dominant (90%). The pressures varied from 10 to 100 MPa. All experiments were conducted at a temperature of 400 °C. The evaluated conditions were chosen to enable comparisons with the experimental data of Seitz and Blencoe (1999).

### 3.2.3. *Dew Line prediction*

The dew lines were evaluated at five different concentrations of 2, 5, 10, 25 and 50% H<sub>2</sub>O. The pressures were chosen to comply with the work done by Patel et al. (1987), and the dew temperatures were then computed dependently.

### 3.2.4. *Rich phases density predictions*

The EoSs were evaluated at four sets of data provided by King et al. (1992), Chiquet et al. (2007), Hebach et al. (2004) and the validated predictions by Bikkina et al. (2011). King et al. (1992) provided only water-rich liquid-phase densities between 6 to 24 MPa at three temperatures of 15, 20, and 25 °C co-existing with CO<sub>2</sub>-rich liquid phase. The 15 and 25 °C test sets were chosen for the analysis here. Chiquet et al. (2007) provided densities of both water-rich and CO<sub>2</sub>-rich phases when CO<sub>2</sub> were at supercritical conditions. The set of data covered pressures from 5 to 45 MPa, and the selected sets of temperatures were about 35, 50, 90 and 110 °C. The work done by Hebach et al. (2004) were used for comparisons of water-rich liquid phase densities co-existing with CO<sub>2</sub>-rich gas phase. The selected cases were at temperatures of 19, 29, 39 and 49 °C and pressures less than 5 MPa to ensure a gaseous CO<sub>2</sub>-rich phase. The results of Chiquet et al. (2007) and Hebach et al. (2004) were measured at temperatures slightly around the listed values, although precisely fixed for each point. The simulations conducted here uses the exact measurements temperature nodes. The group of Bikkina et al. (2011) provided validated predictions that cover the missing rich phases density of CO<sub>2</sub>-rich liquid and the co-existing water-rich liquid densities and CO<sub>2</sub>-rich gaseous phase. The selected Liquid Liquid Equilibrium (LLE) data of Bikkina et al. (2011) went over pressures between 8 to 21 MPa at one temperature of 25 °C. Finally, the Vapour Liquid Equilibrium (VLE) data went over pressures between 1 to 6 MPa at four temperatures of 25, 40, 50 and 60 °C.

### 3.2.5. *Solubilities*

The accuracies of the SPUNG and SRKs EoSs in predicting the mutual solubilities of CO<sub>2</sub> and H<sub>2</sub>O were validated against experimental data. Pappa et al. (2009) reviewed the experimental data of CO<sub>2</sub> -water system solubilities and recommended six sets of mutual solubilities experimental data for model regression and validations. The six sets were of Takenouchi and Kennedy (1964), Wiebe (1941), Bamberger et al. (2000), Valtz et al. (2004), Mueller et al. (1988), and King et al. (1992). For this work we replaced King et al. (1992) with the recent Hou et al. (2013), which covered the available intermediate pressures data

available in literature at various temperatures and filled in some gaps. The latter work also validated the new data against the available literature data and showed very good match. Takenouchi and Kennedy (1964) provided solubilities at very high pressures, ranging between 10 and 70 MPa at a temperature of 110 °C. The study of Wiebe (1941) covered pressures between 1 to 70 MPa and temperatures between 25 and 100 °C. Here, however, we present only the results at 50 and 75 °C since mutual solubilities were provided only at these two temperatures. Bamberger et al. (2000) provided data over moderate pressures between 4 and 14 MPa at three temperatures of 50, 60 and 80 °C. The set of experimental data by Valtz et al. (2004) covered very low pressures at three temperatures of approximately 5, 25 and 45 °C. For the low temperature of 5 °C, the pressures ranged approximately between 0.5 and 1 MPa, the range went wider as the temperature increased reaching approximately the range of 0.1 to 7 MPa at the temperature of 45 °C. The set by Mueller et al. (1988) provided solubility data at low pressures and high temperatures between 100 and 200 °C. The last set of experimental data by Hou et al. (2013) covered a wide range of data for pressures around 1 to 17.5 MPa, and over temperature range of 25 to 175 °C.

### 3.2.6. Sensitivity to the interaction parameter $K_{ij}$

The SPUNG EoS uses the SRK EoS for computing the shape factors and SRK EoS uses the symmetric interaction parameter  $K_{ij}$  for computing the mixing coefficients. Therefore, we performed a simple sensitivity study on the effects of tuning  $K_{ij}$  on the results. The tuning was done by matching the CO<sub>2</sub> solubilities as good as possible and letting the density and H<sub>2</sub>O solubility be computed accordingly. The results of the tuning for the evaluated cases are plotted and labeled as *SPUNG-Reg  $K_{ij}$*  in the following.

### 3.2.7. Reference fluid sensitivity

An examination was conducted for the impact of using other reference fluids on the density predictions of the water-rich liquid phase. N<sub>2</sub>, O<sub>2</sub>, ethane (C1), methane (C2), iso-butane (IC4) and normal-butane (NC4) were used as a reference fluid for this study as an alternative to the originally used propane (C3). The aim of this part of the analysis was to find a criterion of selection or to search for the proper reference fluid to model CO<sub>2</sub>-water mixtures.

### 3.3. Error definition

The errors of an EoS is measured here by the Relative Error (RE) and the Average of Absolute Deviation (AAD) defined for an arbitrary variable  $C$  as

$$RE(C_r, \%) = \frac{|C_{s,r} - C_{\text{exp},r}|}{C_{\text{exp},r}} \times 100 \quad (1)$$

and

$$AAD(C, \%) = \frac{100}{N} \times \sum_{r=1}^N \frac{|C_{s,r} - C_{\text{exp},r}|}{C_{\text{exp},r}} \quad (2)$$

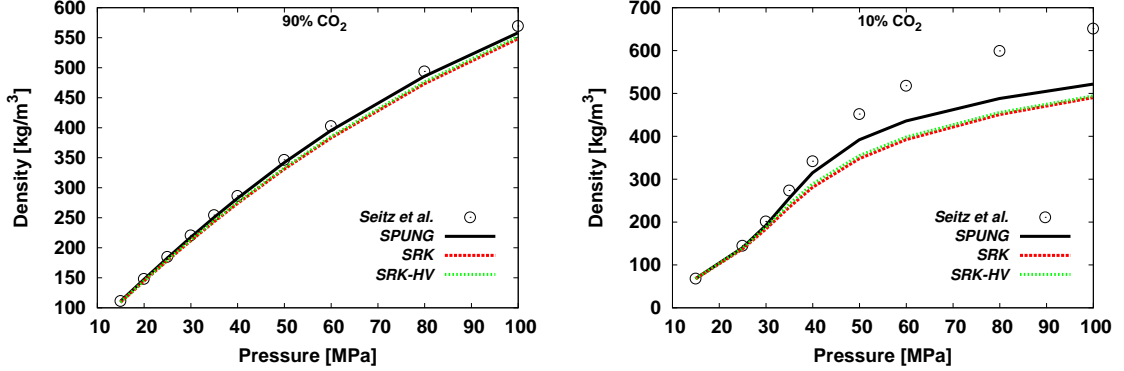


Figure 1: Density computations in comparison with experimental data of Seitz and Blencoe (1999), over pressures up to 100 MPa at 90% and 10% CO<sub>2</sub> and a temperature of 400 °C

Here,  $N$  is the total number of points, subscripts  $s$  and  $exp$  refer to simulation data and experimental data, respectively, and  $r$  is a point index.

## 4. Results

### 4.1. Single phase density, high pressures

The comparisons between the three EoSs for the high-pressure data sets of Seitz and Blencoe (1999) are plotted in Figs. 1 and 2. The former shows the density changes over pressures for various CO<sub>2</sub> content in the mixture, while the latter is an interesting re-plot of the data as density change over molar fraction of CO<sub>2</sub> for the various pressures.

### 4.2. Single phase density, low to moderate pressures

A comparison between the three EoSs for the low pressure data sets of Patel and Eubank (1988) are plotted in Fig. 3 for 2% H<sub>2</sub>O and a temperature of 225 °C. Results for the three other temperatures of 50, 100 and 200 °C were virtually similar to those shown for 225 °C. Furthermore, the models were compared to experimental data for the concentrations of 10% H<sub>2</sub>O at temperatures 71, 100, 200 and 225 °C, for 25% H<sub>2</sub>O at 100, 200 and 225 °C and for 50% H<sub>2</sub>O at 125, 200 and 225 °C. For all these series of density variation with pressure, the errors were tiny and showed no significant dependency on temperature.

### 4.3. Dew line

The dew line predictions of SPUNG, SRK and SRK-HV are shown in Fig. 4 for three of the datasets from Patel et al. (1987). For another two datasets, 25% and 50% H<sub>2</sub>O, the computed results matched the experimental data as good as for 10%, or better. At the compositions of 10% and 25% H<sub>2</sub>O all



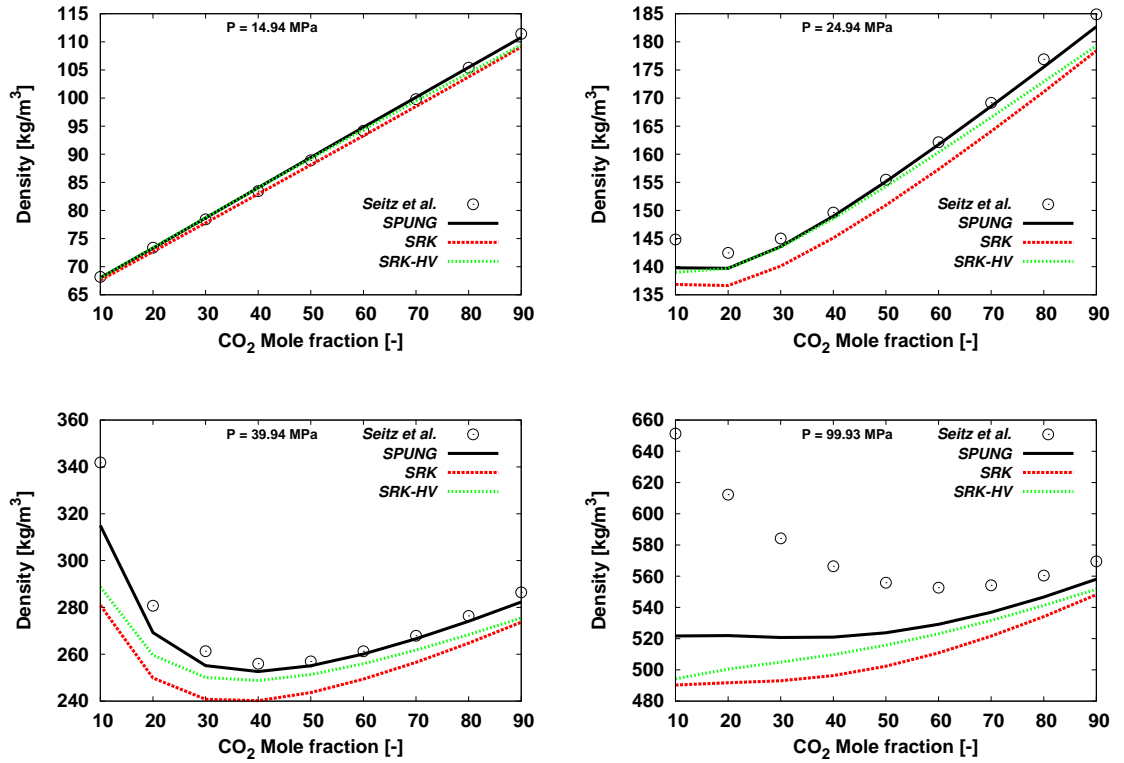


Figure 2: Density computations in comparison with experimental data of Seitz and Blencoe (1999), over mole fractions of CO<sub>2</sub> at different pressures and a temperature of 400 °C

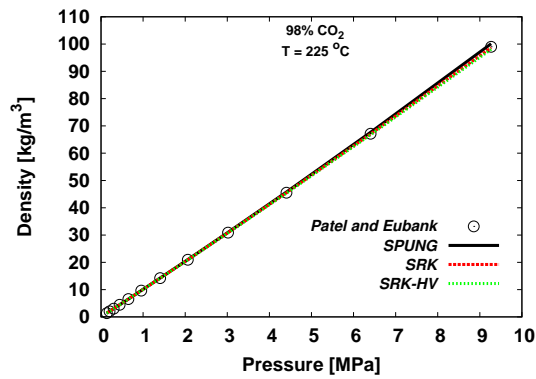


Figure 3: Gas phase density calculations in comparison with experimental data of Patel and Eubank (1988) at 2% H<sub>2</sub>O and a temperature of 225 °C

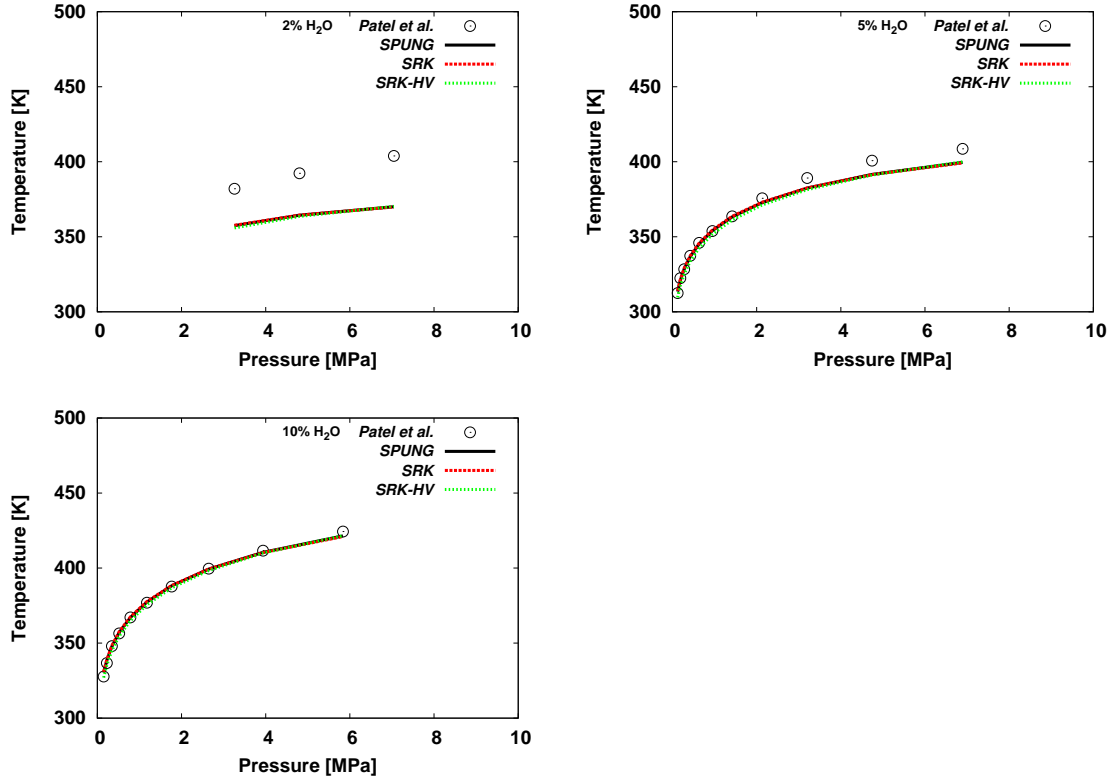


Figure 4: Dew line temperature predictions in comparison with experimental data of Patel et al. (1987)

the EoSs predicted pseudo critical pressures lower than the highest pressure of the experiments. SRK-HV predicted 8.045 and 9.394 MPa pseudo critical pressures, respectively, and SRK predicted 8.161 and 9.55 MPa, respectively, for the two compositions. Since SPUNG uses the SRK algorithm to calculate for the pseudo critical quantities, it predicted the same pseudo critical pressures as of SRK. Hence, the highest pressure point of the 10% H<sub>2</sub>O content from Patel et al. (1987) was not plotted. As seen in Fig. 4, the results showed an improvement of the EoSs predictions as the water content increased.

#### 4.4. Rich phases density prediction

##### 4.4.1. CO<sub>2</sub>-rich phases

The densities of the Supercritical Liquid Equilibrium (SGLE) of the CO<sub>2</sub>-rich phase co-existing with a liquid water-rich phase were modeled, and the results are presented in comparison with the experimental data of Chiquet et al. (2007)

Table 1: AAD [%] of the CO<sub>2</sub>-rich phases density calculations for the CO<sub>2</sub> -water system at different CO<sub>2</sub> co-existing phases and temperatures

Data sets	Phase Equilibrium	Temperature [°C]	SPUNG	SRK	SRK-HV	SPUNG-Reg $K_{ij}$
<i>Chiquet et al. (2007)</i>	SGLE	35	1.8	8.3	8.7	4.5
		50	0.7	6.6	7.4	4.5
		90	4.5	6.9	8.2	7.1
		110	5.3	7.1	7.9	7.4
		90 <sup>a</sup>	2.7	5.3	6.7	5.7
		110 <sup>b</sup>	2.0	5.9	7.3	3.7
<i>Bikkina et al. (2011)<sup>c</sup></i>	LLE	25	2.0	9.2	9.7	8.3
<i>Bikkina et al. (2011)<sup>c</sup></i>	VLE		0.6	0.3	0.3	0.7

<sup>a</sup> without the anomalous point 7.5 MPa, <sup>b</sup> without the anomalous point 25 MPa, <sup>c</sup> predicted data

in Fig. 5. The AAD of the supercritical CO<sub>2</sub>-rich phase density predictions are presented in Table 1.

From Fig. 5, two experimental points seemed to be anomalous: at 110 °C, 25 MPa and 90 °C, 7 MPa. These points deviated from the trend of each dataset, and the model errors jumped significantly. For the discussion, the AADs were recalculated without these two points in Table 1.

Model computations of the liquid CO<sub>2</sub>-rich phase were compared with the predicted data of Bikkina et al. (2011) in Fig. 6. The corresponding AADs are included in Table 1.

For the gaseous CO<sub>2</sub>-rich phase predictions (results are not plotted), errors compared with the data of Bikkina et al. (2011) were very small with all models and very similar to the single gas phase results. However, the values of the binary interaction parameter  $K_{ij}$  used to get the proper CO<sub>2</sub> solubility decreased with increasing temperature. The used values were  $-1.44$ ,  $-0.130$ ,  $-0.115$  and  $-0.107$ , respectively, for 25, 40, 50 and 60 °C. The AADs for the entire used dataset are included in Table 1. For the gaseous phase, the AADs reported were temperature averaged.

#### 4.4.2. Water rich liquid density prediction

The density predictions of the liquid water-rich phase co-existing with a supercritical CO<sub>2</sub>-rich phase at a temperature of 35 °C are presented in comparison with the experimental data of Chiquet et al. (2007) in Fig. 7. The results for the temperatures of 50, 90 and 110 °C were very similar in trend. However, the  $K_{ij}$  values used to get the proper CO<sub>2</sub> solubility decreased with temperature increase, where the used values were  $-0.132$ ,  $-0.118$ ,  $-0.068$ ,  $-0.045$ , respectively, for the temperatures from 35 to 110 °C.

The results of the density predictions of the liquid water-rich phase co-existing with liquid CO<sub>2</sub>-rich phase at a temperature of 25 °C are plotted in Fig. 8. The results of 15 °C behaved very similar to those at 25 °C and are not shown. The used  $K_{ij}$  values were  $-0.15$  and  $-0.14$ , respectively.

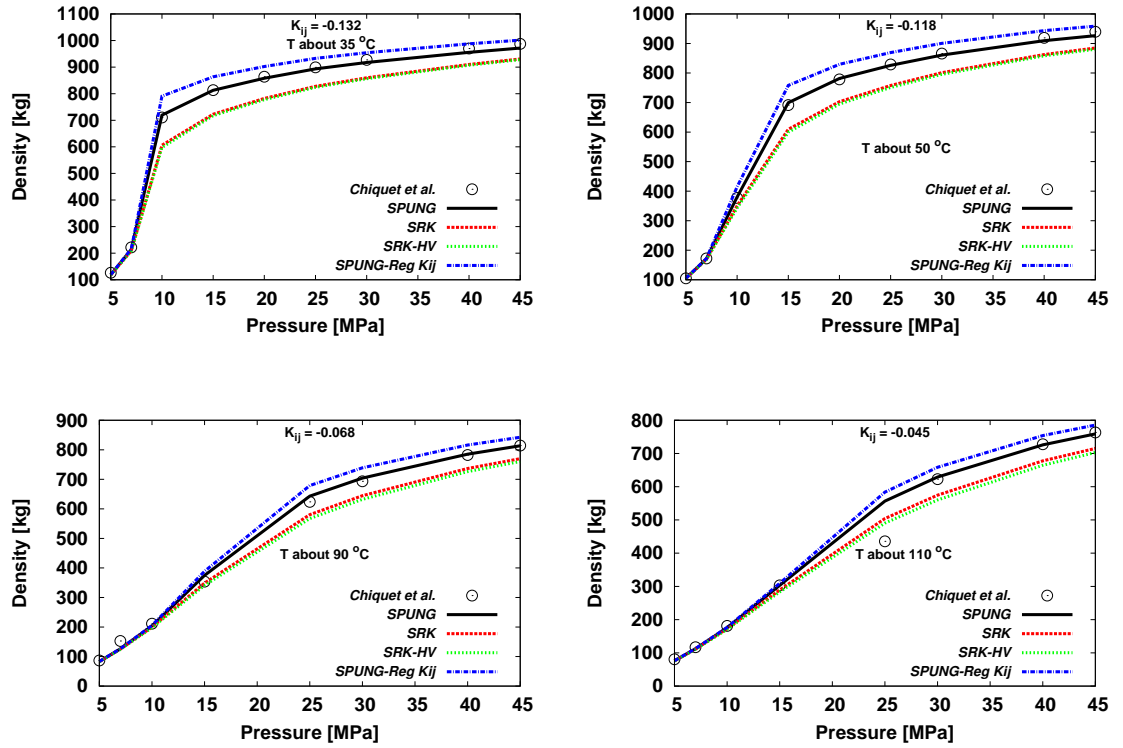


Figure 5: SGLE CO<sub>2</sub>-rich phase density predictions at temperatures about 35, 50, 90 and 110 °C in comparison with Chiquet et al. (2007) experimental data

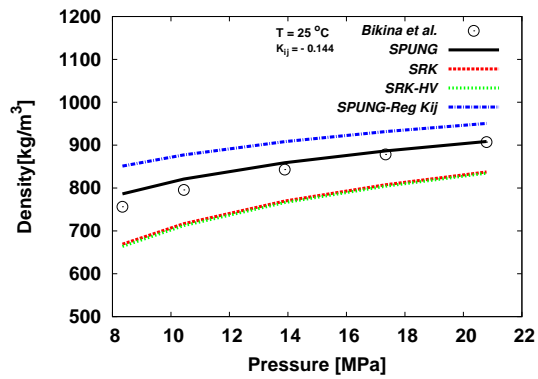


Figure 6: LLE CO<sub>2</sub>-rich phase densities prediction at a temperature of 25 °C in comparison to predictions of Bikina et al. (2011)

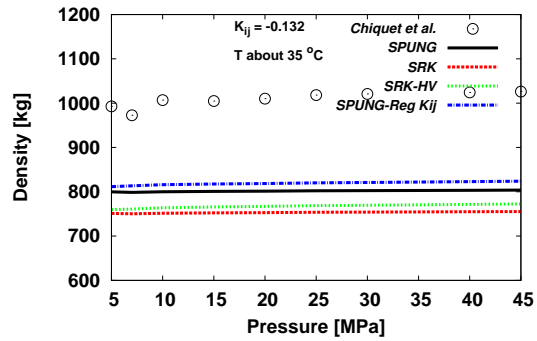


Figure 7: Densities of the liquid water-rich phase co-existing with a supercritical CO<sub>2</sub>-rich phase at a temperature about 35 °C in comparison with Chiquet et al. (2007) experimental data

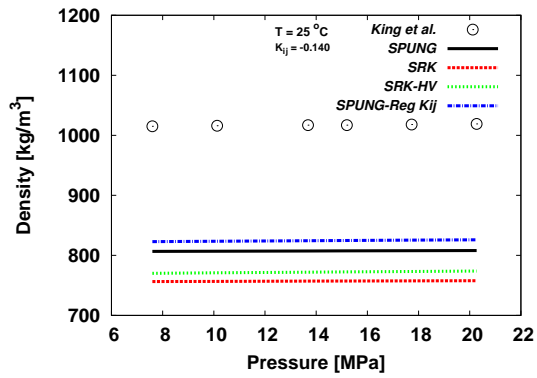


Figure 8: LLE water-rich phase density predictions at a temperature 25 °C in comparison with King et al. (1992) experimental data

Table 2: AAD [%] of Water rich liquid phase densities averaged over the temperatures of each evaluated set of data

Data sets	Phase Equilibrium	SPUNG	SRK	SRK-HV	SPUNG-Reg $K_{ij}$
<i>King et al. (1992)</i>	LLE	20.5	25.5	23.9	18.7
<i>Chiquet et al. (2007)</i>	SGLE	21.6	26.0	25.0	20.3
<i>Hebach et al. (2004)</i>	VLE	20.1	24.9	24.5	19.4

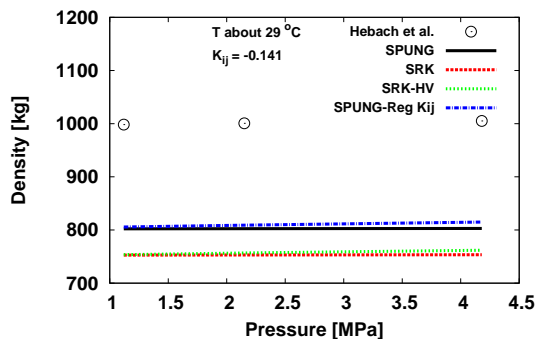


Figure 9: Densities of the liquid water-rich phase co-existing with a gaseous  $\text{CO}_2$ -rich phase at a temperature about  $29^\circ\text{C}$  in comparison with Hebach et al. (2004) experimental data

Density predictions of the liquid water-rich phase co-existing with gaseous  $\text{CO}_2$ -rich phase at a temperature of  $29^\circ\text{C}$  are plotted in Fig. 9. The results of 19, 39 and  $49^\circ\text{C}$  were very similar. The used  $K_{ij}$  values were  $-0.154$ ,  $-0.141$ ,  $-0.129$  and  $-0.118$ , respectively. Table 2 contains a summary of the results in terms of temperature averaged AADs.

As mentioned in Sect. 2.4.2, propane was chosen as the reference fluid in the SPUNG EoS. Table 3 shows the results from the reference fluid sensitivity study. The AADs is averaged over temperature. The C3 (propane) results were the same as shown for SPUNG in Fig. 9. Only one temperature out of the evaluated four was presented due to similarity in trends and uniformity of the errors. The curves of the REs showed almost equal slopes for the evaluated reference fluids (not shown). The difference in AADs is around 7% between using  $\text{N}_2$  and  $\text{NC}_4$  as a reference fluid.

Table 3: AAD [%] of densities averaged over the temperatures of the comparison with Hebach et al. (2004) data using various reference fluids

Data set	Phase Equilibrium	$\text{N}_2$	$\text{O}_2$	C1	C2	C3	$\text{NC}_4$
<i>Hebach et al. (2004)</i>	VLE	25.7	24.8	24.8	21.8	20.1	18.9

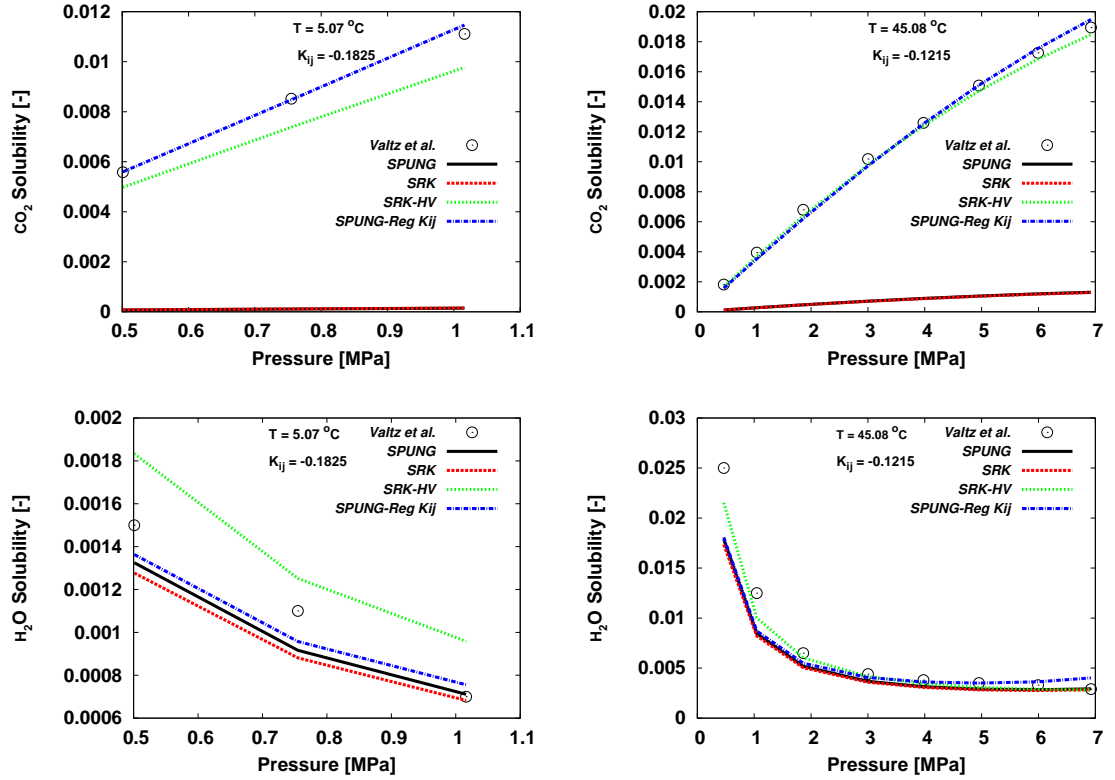


Figure 10: CO<sub>2</sub> and H<sub>2</sub>O solubilities over low pressures and at temperatures of about 5 and 45 °C in comparison with Valtz et al. (2004) experimental data

#### 4.5. Solubilities

The behavior of SPUNG, SRK and SRK-HV at low pressures and low temperatures were evaluated, and results in comparison to the work conducted by Valtz et al. (2004) are plotted in Fig. 10. The solubilities at 25.13 °C were also evaluated towards experimental data with results comparable to those shown. The AADs are presented in Table 4.

The results in comparison to the work conducted by Mueller et al. (1988) are plotted in Fig. 11. The results of the intermediate temperatures are not plotted, since the plotted ones are sufficient to show the trend. The AADs at 100, 140 and 200 °C are presented in Table 5.

The solubilities over moderate pressures were predicted by the three models at the temperatures of 50, 60 and 80 °C, which were chosen in consistency to the experimental work of Bamberger et al. (2000). Results are plotted in Figs. 12 and 13. The CO<sub>2</sub> solubility results at 50 and 60 °C turned out very similar to those shown at 80 °C. For H<sub>2</sub>O the 50 °C results were similar to those

Table 4: AAD [%] of the solubility of CO<sub>2</sub> and H<sub>2</sub>O in comparison to Valtz et al. (2004)

Component	Temperature [°C]	SPUNG	SRK	SRK-HV	SPUNG-Reg $K_{ij}$
CO <sub>2</sub>	5.07	98.7	98.6	12.0	2.7
	25.13	96.8	96.7	8.8	5.5
	45.08	93.1	93.2	4.3	5.3
H <sub>2</sub> O	5.07	10.0	12.4	24.3	10.0
	25.13	17.8	20.8	8.3	17.0
	45.08	17.5	18.9	10.4	16.8

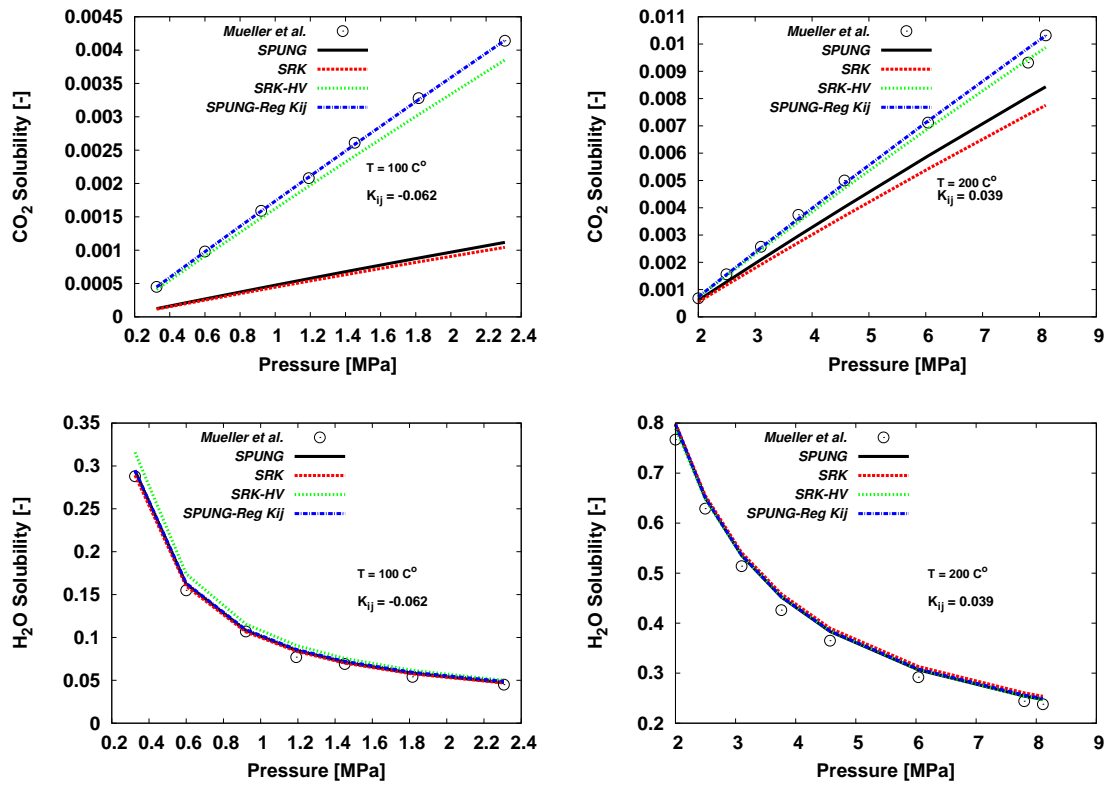


Figure 11: CO<sub>2</sub> and H<sub>2</sub>O solubilities over low pressures and at temperatures of 100 and 200 °C in comparison with Mueller et al. (1988) experimental data



Table 5: AAD [%] of the solubility of CO<sub>2</sub> and H<sub>2</sub>O in comparison to Mueller et al. (1988)

Component	Temperature [°C]	SPUNG	SRK	SRK-HV	SPUNG-Reg $K_{ij}$
CO <sub>2</sub>	100	72.8	74.5	7.2	0.5
	140	49.2	53.2	4.9	1.1
	200	17.0	24.2	4.5	2.8
H <sub>2</sub> O	100	4.9	3.6	11.8	6.3
	140	3.1	3.3	0.9	2.4
	200	4.5	6.1	4.4	4.9

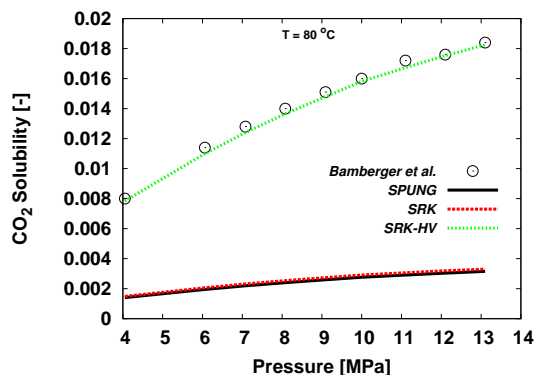


Figure 12: CO<sub>2</sub> solubilities over moderate pressures and at a temperature of 80 °C in comparison with Bamberger et al. (2000) experimental data

shown for 60 °C (Fig. 13), with some better match with experimental data for SRK-HV. The AADs are presented in Table 6. A  $K_{ij}$  sensitivity study was conducted over this set of conditions. The results showed that any improvement of CO<sub>2</sub> solubilities prediction causes a significant increase in the H<sub>2</sub>O solubility prediction errors for SRK and SPUNG EoSs.

The results of the comparisons to Hou et al. (2013) are plotted in Fig. 14. Graphs are included only at the lowest and highest temperatures, as these are sufficient to illustrate the trend.

The comparison with Wiebe (1941) is presented in Table 7. There were no compromise found for SPUNG EoS by tuning  $K_{ij}$ . The results showed that any improvement of CO<sub>2</sub> solubilities prediction causes a significant increase in the H<sub>2</sub>O solubility prediction errors for SRK, and SPUNG EoSs.

The evaluated EoSs were used to predict the mutual solubilities of CO<sub>2</sub> and H<sub>2</sub>O at very high pressures. The conditions were chosen in compliance with the work of Takenouchi and Kennedy (1964). Pressures from around 10 to 70 MPa were used for predictions at a temperature of 110 °C. The results are plotted in comparison with experimental data in Fig. 15. The errors are described in

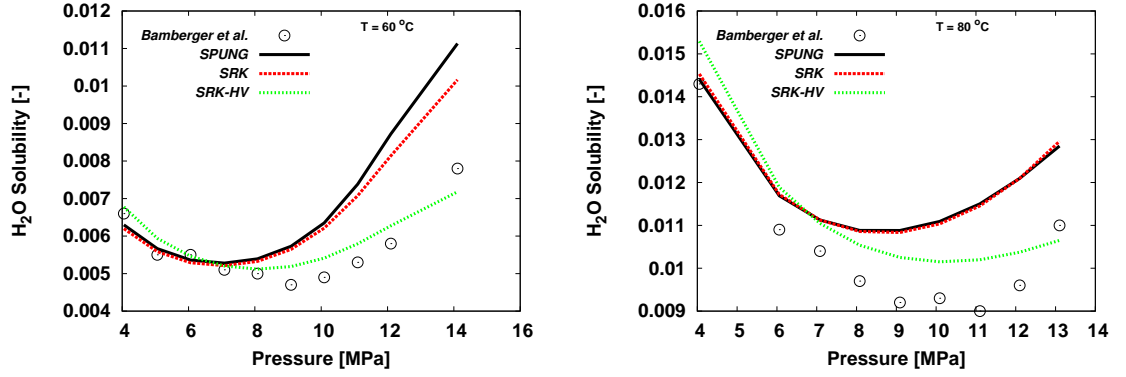


Figure 13: H<sub>2</sub>O solubilities over moderate pressures and at temperatures of 60 and 80 °C in comparison with Bamberger et al. (2000) experimental data

Table 6: AAD [%] of the solubility of CO<sub>2</sub> and H<sub>2</sub>O at different temperatures in comparison with Bamberger et al. (2000)

Component	Temperature [°C]	SPUNG	SRK	SRK-HV
CO <sub>2</sub>	50	91.8	91.7	2.1
	60	88.9	89.2	1.3
	80	82.8	81.9	2.0
H <sub>2</sub> O	50	18.2	24.9	4.2
	60	20.5	17.1	6.2
	80	15.0	15.1	8.4

Table 7: AAD [%] of the solubility of CO<sub>2</sub> and H<sub>2</sub>O at different temperatures in comparison with Wiebe (1941)

Component	Temperature [°C]	SPUNG	SRK	SRK-HV
CO <sub>2</sub>	50	92.0	91.8	6.9
	75	84.3	84.3	2.01
H <sub>2</sub> O	50	167.9	157.9	33.1
	75	148.5	137.2	35.5

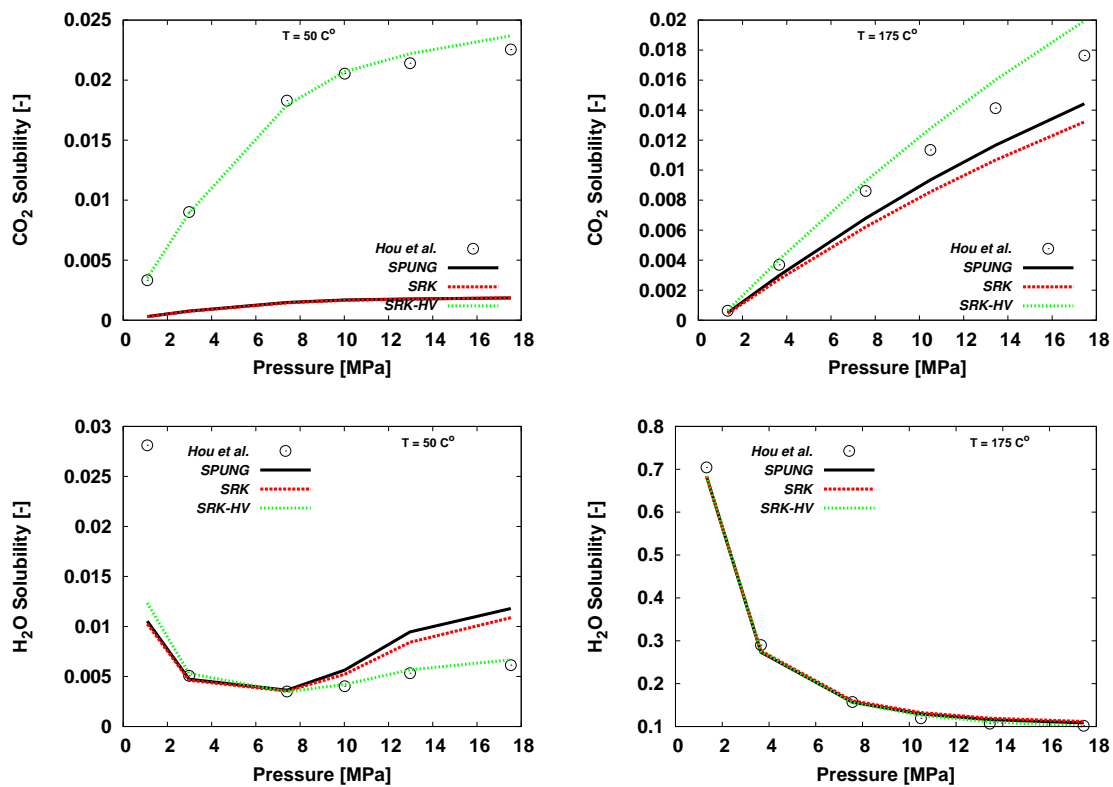


Figure 14: CO<sub>2</sub> and H<sub>2</sub>O solubilities computations in comparison with experimental data of Hou et al. (2013)

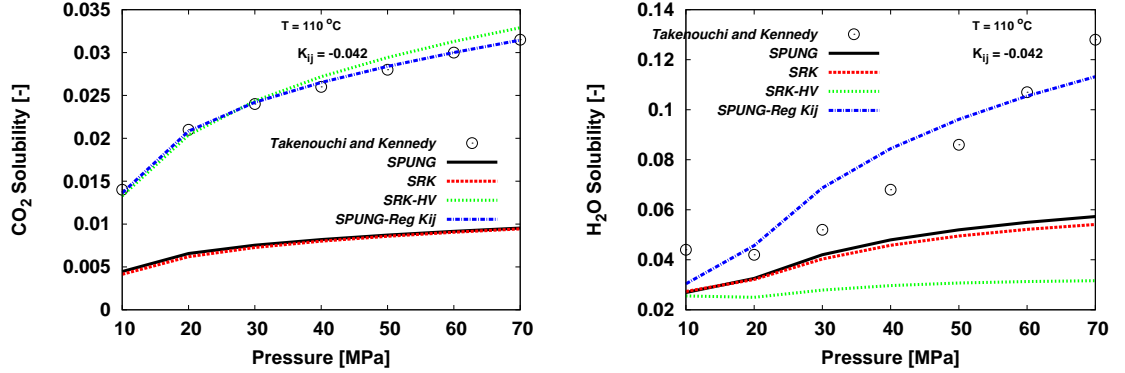


Figure 15: CO<sub>2</sub> and H<sub>2</sub>O solubilities over very high pressures and a temperature of 110 °C in comparison with Takenouchi and Kennedy (1964) experimental data

Table 8: AAD [%] of solubilities in comparison with Takenouchi and Kennedy (1964)

Component	SPUNG	SRK	SRK-HV	SPUNG-Reg $K_{ij}$
CO <sub>2</sub>	68.9	69.9	4.0	1.1
H <sub>2</sub> O	36.2	38.3	56.5	17.3

terms of AADs and presented in Table 8.

## 5. Discussion

### 5.1. Single phase and dew line prediction

Regarding dew line predictions or saturation conditions, Fig. 4 and results for higher H<sub>2</sub>O content showed that all the three EoSs were behaving well in computing the saturation line. An exception is the case of 2% H<sub>2</sub>O, which seems to be challenging for all the models. The simulations also showed that all the tested models predicted a low pseudo-critical pressure for the cases of 10% and 25% H<sub>2</sub>O.

For low-pressure density calculations presented in Fig. 3, the densities comparison showed that the errors over the investigated intervals were on average very small for all the tested EoSs. Nevertheless, looking closely at the errors behavior, it was observed that the errors increased as the pressure increased and the mixture gradually departed from ideal behavior. The important observation was that the errors of SRK and SRK-HV grew steeply compared to that of the SPUNG EoS as the pressure went above a certain threshold in most of the cases. This behavior resulted in REs of SRK and SRK-HV that were multiples of that of SPUNG at the upper bound of the tested (pressure and H<sub>2</sub>O) intervals.

From the results of the high-pressure single-phase density calculations Figs. 1 and 2 the comparisons showed an increase of the errors as the pressure increased

and as the H<sub>2</sub>O content increased. While at low pressure the increase and the relative errors value were small, the errors jumped to an order of magnitude higher at the combination of the upper-bounds of both intervals.

The figures also show clearly that the errors behavior of SPUNG is much better than that of SRK and SRK-HV and, considering the computational expenses study by Wilhelmsen et al. (2012), it can be concluded that it is a good compromise between sophisticated and cubic EoSs.

Although the errors of SPUNG reached 20% at the extreme of the investigated conditions, the method has a possible high potential for improvement via using other reference fluids, while the SRKs do not have the same potential. Further research can evaluate this potential.

The inaccuracies of the used cubic EoSs are due to the simple structure of the models, which have very few parameters to tune. A study similar to the presented work but for other mixtures was made by Li and Yan (2009), who reported the same inaccuracies using SRK and other cubic equations for mixtures of CH<sub>4</sub>, H<sub>2</sub>S, N<sub>2</sub> and Ar. Furthermore, Li et al. (2011) reviewed several studies testing cubic equations for gas and liquid density predictions for other mixtures. In our investigation, the errors reached approximately 25% at the extreme conditions using cubic EoSs. This was higher than in the studies of other mixtures, emphasizing how challenging this particular mixture is for cubic EoSs compared to the other mixtures. In addition, this showed the need for a more predictive concept when dealing with CO<sub>2</sub> -water mixture.

On the other hand, the SPUNG EoS superiority in density computations was inherited from the use of the 32-parameter MBWR reference equation, which is very accurate for propane. However, the errors of the SPUNG EoS came from the incapability of propane to achieve the high density of the CO<sub>2</sub> - water liquid phase.

## *5.2. Rich phases density prediction*

### *5.2.1. CO<sub>2</sub>-rich gas phase*

The results of CO<sub>2</sub>-rich gas phase showed that the accuracy of all the evaluated EoSs was very good. A high accuracy for gas phase density using cubic EoSs was reported in many studies that were listed in the review article of Li et al. (2011). Furthermore, the solubility of the H<sub>2</sub>O into CO<sub>2</sub>-rich gas phase was too small to cause a challenge in modelling, as reported by Hebach et al. (2004).

### *5.2.2. CO<sub>2</sub>-rich liquid phase*

The results in Fig. 6 show that at the low pressure side, the errors of the SRKs were around 12% compared to 4% of SPUNG EoS. As the pressure went higher, the errors of all the EoSs became lower. However, the errors of SPUNG dropped to around 0.17 while the SRKs errors remained high at approximately 8%. This behavior was not revealed by the AADs in Table 1, which average the REs over the predicted interval to approximately 2% of SPUNG and 9% for SRKs.

### 5.2.3. CO<sub>2</sub>-rich supercritical phase

As noted in Sect. 4.4.1, two measurement points, (7 MPa, about 90 °C) and (25 MPa, about 110 °C), in Fig. 5 seemed to be anomalous. The deviation is seen for all the three evaluated models. Since the models are based on different theories, the anomaly suggested a measurement error. Alternatively, there might be a feature that is not captured by any of the models. At the two low evaluated temperatures, the predictions of SPUNG EoS were substantially better than those of the SRKs, especially in capturing the steep change in density over the pressures between 5 and 10 MPa at about 35 °C and between 5 and 15 MPa at about 50 °C, as observed in Fig. 5. In these two cases, the errors of SRKs jumped to around 15 and 13%, while the SPUNG errors were below 1.3%. Furthermore, the errors of the SRKs were reduced gradually as the density to pressure curve started saturating, while the errors of SPUNG remained low over the entire interval. Table 1 flattened out this behavior to AADs, which in their turns showed the high accuracy of SPUNG EoS prediction compared to SRKs EoSs. At the two evaluated high temperatures, the two mentioned points gave exceptional peaks in the error. Apart from this, the errors were similar to those of the lower temperatures, although the amplitude was much lower, and the RE distribution in general had a more flattened profile as the density increase with pressure was much more gradual at high temperatures than at the low temperature cases. The behavior of SPUNG remained superior, which can also be observed in the results summarized in terms of AADs in Table 1.

### 5.2.4. Water-rich liquid phase

The results in Fig. 7 and the results for other temperatures (not presented) showed that the errors of all the evaluated EoSs were considerable especially when compared to the results of Tsivintzelis et al. (2011) and Diamantonis and Economou (2012) for CPA and PC-SAFT, respectively. The errors were not very sensitive to temperature and pressure. This caused the RE profiles to be rather flat and made the AAD a very representative measure.

The capture of the temperature dependency was good. Furthermore, the insensitivity to both the pressure and co-existing phase was virtual as it was due to the incapability of all the EoSs to capture the high liquid-water density. This is regardless of how well the models capture the CO<sub>2</sub> solubility effect due to the increase in pressure. However, looking carefully to the SRK-HV slope in density-pressure behavior and the rather horizontal REs compared to the other evaluated EoSs, it could be observed that only SRK-HV captured the effect of CO<sub>2</sub> solubility as a function of pressure properly due to the superior prediction of CO<sub>2</sub> solubility using SRK-HV, which will be discussed below. This observation was supported by comparing the predictions of pure water at the same pressures and the one with CO<sub>2</sub> solubility, where SPUNG and SRK showed almost no difference in density predictions, whereas SRK-HV predicted the density difference accurately. This observation was not very clear from the first glance at the graphs, since the difference it made to capture the CO<sub>2</sub> solubility properly was of 1.5%, while the errors were above 20% for all the evaluated

EoSs. Although SRK-HV predicted the deviation part correctly, SPUNG density prediction was superior to both SRKs, with a potential of improvement by using other reference fluids.

This discussion applies to the results of water-rich liquid phase co-existing with liquid and gaseous CO<sub>2</sub>-rich phases in Figs. 8-9. Nevertheless, the difference in the case of water-rich liquid co-existing with gaseous CO<sub>2</sub>-rich phase was that the measured density had a slight increase in with increased pressure, Fig. 9. This is due to the interfacial tension as reported by Hebach et al. (2004). The solubility was captured very good by SRK-HV. This can be observed in the inclination of the SRK-HV curve, which has the very similar slope as of the experimental data in the density results of Fig. 9.

An observation from Fig. 7 was that the point of 7 MPa seemed to deviate from the trend of the remaining points. Unless this was just an inaccuracy, the phenomenon was not captured by the models.

### 5.3. Solubilities

The results in Fig. 15 and Table 8 show that the predictability of SRK-HV for the solubility of CO<sub>2</sub> in water was much better than those of SRK and SPUNG EoSs and of low errors. The predictions by SPUNG and SRK were poor. On the other hand, the prediction of the H<sub>2</sub>O solubility by SRK-HV was much worse than that of SRK and SPUNG, where all the models were inaccurate. Since SRK uses a symmetric interaction parameter  $K_{ij}$  between CO<sub>2</sub> and H<sub>2</sub>O in the van der Waals geometric mean-based mixing rules, it was expected that the SRK predictability of one of the mutual solubilities will be that low due to the polar nature of the mixture. The results suggested that SPUNG EoS inherits this impotence from SRK since it uses SRK to compute the shape factors.

The comparison with Wiebe (1941) at high pressures and low temperatures showed very low predictability of all the EoSs with very deviating results using both SPUNG and SRK EoSs. There were no improvement for SPUNG and SRK EoSs achieved by the regression.

The solubilities at moderate pressures, chosen in consistency to the experimental work conducted by Bamberger et al. (2000), were predicted by the three models. The results in Figs. 12 and 13 and AADs in Table 6 show that the errors in predicting CO<sub>2</sub> solubility became more severe than those at very high pressures for SPUNG and SRK EoSs. This highlighted the superior behavior of SRK-HV even more. Furthermore, the errors of SPUNG and SRK were reduced as the temperature increased, which suggested a need for correlating the interaction parameter  $K_{ij}$  to temperature in addition to a more general mixing rule. This analysis was confirmed by the  $K_{ij}$  sensitivity study conducted here (see Sect. 5.4). The predictability of SPUNG and SRK improved for H<sub>2</sub>O solubilities, while that of SRK-HV improved for both mutual solubilities and behaved much better than those of SPUNG and SRK.

The comparison to the set of data of Hou et al. (2013) showed the same behavior as of the one with Bamberger et al. (2000) at similar temperatures. However, the errors of all the EoSs were reduced significantly as the temperature increased.

For low pressures, the results in Fig. 10 and the AADs in Table 4 showed the same trend in comparison with the experimental data of Bamberger et al. (2000), except that SRK-HV did not behave equally well.

For the low pressures and high temperatures in Fig. 11, the results showed good and improving predictability as the temperature increased.

In general, the predictions of all the EoSs improved with temperatures increase at all pressures.

#### 5.4. Effects of the interaction parameter $K_{ij}$

Since the impotence of SPUNG in predicting the solubilities of CO<sub>2</sub> - water was thought to be inherited from SRK due to the use of the symmetric interaction parameter  $K_{ij}$ , a simple sensitivity study on the effects of tuning  $K_{ij}$  was conducted as explained in Sec. 3.2.6. The results of the tuning for the evaluated very high and low-pressure cases are plotted in Figs. 10, 11 and 15. The AADs are presented in Tables 4, 5 and 8. These results showed that at these conditions there existed a  $K_{ij}$  that could improve the mutual solubilities together, and which compromised the errors better than SRK-HV for the very high pressure, and for the low pressures-high temperatures cases. This implies a potential improvement by regression. Unfortunately, this behavior did not hold for the moderate pressures, and high pressures low temperatures cases. There, the CO<sub>2</sub> solubility errors could be improved, but causing the errors of H<sub>2</sub>O solubility to jump high with the expected counter effects due to the use of the geometric mean mixing rules. This behavior shows clearly that SPUNG EoS solubility prediction is limited to that of the EoS used to compute the shape factors.

The tuning of  $K_{ij}$  influenced the rich-phase density predictions mainly through solubility. This was because  $K_{ij}$  influences only the energy parameter  $a$ , and not the co-volume parameter  $b$  in the cubic EoS formulations. In order to give a sense of the impact of each of the mutual solubility on the density predictions for CO<sub>2</sub> -water system, a test was conducted at the conditions used for the density analysis discussed here. Since Chiquet et al. (2007), Hebach et al. (2004) and Bikkina et al. (2011) have not provided solubilities, the tuning was done by matching SRK-HV CO<sub>2</sub> solubility as good as possible. This was thought to be a valid step because SRK-HV showed a significantly good prediction of CO<sub>2</sub> solubility at similar pressure ranges in the work of Austegard et al. (2006). In addition, the comparison here, Figs. 12 - 13 and Table 6, supported the same claim. The density predictions of the  $K_{ij}$  tuned SPUNG EoS were included in the figures and summarized in the tables of the density predictions study. The results showed minor improvement on H<sub>2</sub>O-rich phase density and major dis-improvement on CO<sub>2</sub>-rich phase density due to the conjugate dis-improvement in H<sub>2</sub>O solubility prediction that was induced from the mixing rule.

It is important also to highlight the observation that the tuning showed a temperature-dependent behavior for  $K_{ij}$  that was almost insensitive to the pressure and co-existing phases.

Since the solubilities are important for deciding co-existing phases, especially for small impurities of H<sub>2</sub>O in CO<sub>2</sub> or vice versa, and since SPUNG EoS has



shown high potential, we are motivated under the guidance of this work for further developing the method to overcome this weakness. A more elaborated mixing rule that shall take into account the polar nature of the system, as well as the temperature dependency shown in this work, is needed.

### 5.5. Reference fluid sensitivity

The results in Table 3 showed clearly that the choice of reference fluid had a significant impact on the properties predictions and, in particular, on density predictions. Furthermore, the trend observed in the results was very interesting, where a heavier reference fluid within the set of hydrocarbons gave better predictions of the density compared to a lighter. Also within the set of O<sub>2</sub> and N<sub>2</sub> the same trend was seen. For all the reference fluids in Table 3, the curves of the REs showed almost equal slopes, which implied a low impact on solubility predictions.

## 6. Conclusions

The three tested EoSs predicted the dew temperature with high quality and precision, but predicted low pseudo critical pressures for two tested data sets.

For single phase, at low pressure gas phase, SPUNG EoS exhibits a better behavior to SRK and SRK-HV cubic EoSs. However, the relative errors are low for all models. The role of SPUNG becomes significant as high pressures are of concern, where the error become considerable.

SPUNG has a superior behavior in predicting the rich phases density of the CO<sub>2</sub>-water system compared to the evaluated cubic EoSs. Although CO<sub>2</sub> solubility prediction of SPUNG is very low at moderate pressures and low temperatures, the impact on density calculations for the H<sub>2</sub>O-rich phase is not pronounceable. Improving the CO<sub>2</sub> solubility on the benefits of that of H<sub>2</sub>O, leads to severe mis-prediction in the density of the CO<sub>2</sub>-rich phase. The impact on the overall density prediction of the system will depend on the feed composition. Therefore, for the cases where water is an impurity, the impact of CO<sub>2</sub> solubility mis-prediction will have much less impact on the overall density prediction.

The effect of varying the reference fluids was investigated, and the errors span between the lightest and the heaviest reference fluid was large. This implies a significant impact of the reference fluid on the properties prediction. Nevertheless, the heaviest evaluated hydrocarbon was not heavy enough to give a significant improvement. However, the observed trend and highlighted criterion of the search for a reference fluid rises the expectations in the SPUNG EoS potential for improving the water-rich phase density prediction, if a proper reference fluid is found, while the cubic EoSs do not have a similar potential.

SRK-HV EoS predicted the mutual solubilities for the binary polar mixture with high accuracy. Nevertheless, it showed much poorer predictability of the density of the CO<sub>2</sub>-water system in general and compared to SPUNG in particular.

SRK EoS with van der Waals mixing rules combines the impotence of both SPUNG and SRK-HV EoS. Therefore, it is not recommend for this system, unless low-pressure gas-phase densities are the only interest.

The study showed that the SPUNG EoS predictability of the mutual solubilities is limited by the EoS used for the computation of the shape factors, which here, was SRK. However, the predictability of the density depends more on the choice of the reference fluid and the reference equation used. Since one of the powerful features of the concept is to allow a free choice of the EoS for the shape factors, the reference fluid, and the reference equation (given that the reference fluid coefficients exist for this reference equation), a promising alternative is to use an asymmetric quadratic mixing rule. It is also possible to use SRK-HV, which showed a very high success for solubility predictions of CO<sub>2</sub>-water system. The work shows that the mixing rule has to have parameters fitted at each temperature for CO<sub>2</sub>-water system.

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